# Performance of different modelling approaches for the simulation of pile driving process and for the estimation of the dynamic foundation-subsoil interaction of a wind turbine

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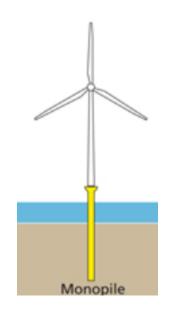






#### **Content**

i. FE model for the simulation of the vibratory pile installation technique in saturated soils



ii. Development of an engineer-oriented model for monopile foundations based on the HCA model

Wind load on Blades
1P load

Resultant 3P load

Wave load

iii. Estimation of the dynamic foundation-subsoil interaction of a wind turbine





## i. FE model for the simulation of the vibratory pile installation technique in saturated soils





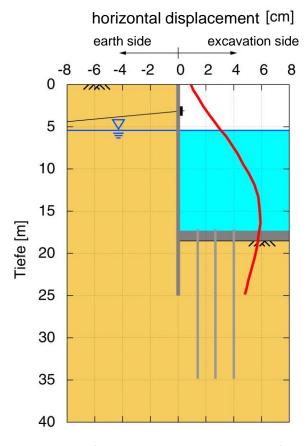
#### **Motivation**

#### Vibratory pile driving in saturated soil

- Deformations in the surrounding soil
- Stress redistributions (Decrease of effective stresses)
- Large displacements of adjacent structures



(Aubram et al., 2015)













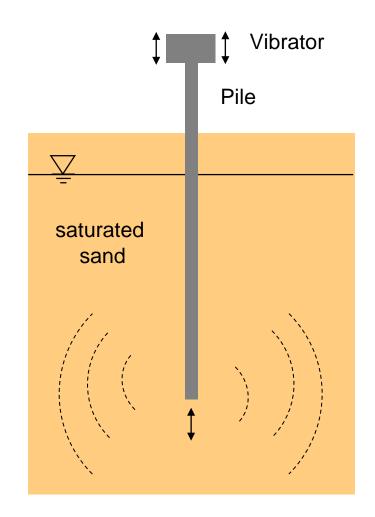
#### **Motivation**

#### Demand:

Development of a reliable modelling approach for the simulation of soil deformation and stress redistribution during vibratory pile installation process in water-saturated soil

#### **Difficulties:**

- Interaction problem with large material deformations
- Dynamic problem in water-saturated soil
- High requirements for suitable constitutive models due to cyclic deformation process
  - Only few studies in the literature
  - No validation with experimental data









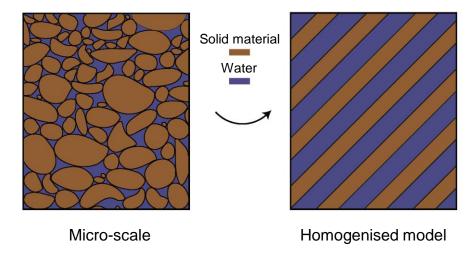
#### Methodology

- Development and implementation of a user-defined element in Abaqus/Standard for the dynamic analysis of fluid-saturated soils
- Establishing of FE models for the simulation of vibratory pile driving
- Quantitative verification of the numerical modelling approaches based on comparisons with model tests for vibratory pile driving
- Investigation of whether the FE models can reproduce the essential aspects of vibratory pile driving





#### **Theory of Porous Media**



- Constituents are present as continuous phases
- Considering as a homogenised continuum
- u-p-approximation for a two-phase medium (Zienkiewicz, 1983)
- Axisymmetric 2D-formulation







### Strong form of the equations (u-p-Approximation)

• Equations of motion

$$\operatorname{div} \boldsymbol{\sigma} - \operatorname{grad} p + \varrho \, \mathbf{g} = \varrho \ddot{\mathbf{u}}$$
$$-\operatorname{grad} p + \varrho_f \mathbf{g} - \frac{1}{k} \, \varrho_f g \mathbf{w} = \varrho_f \ddot{\mathbf{u}}$$

• Constitutive equation for the pore pressure

$$\dot{p} = -\frac{K_f}{n} \operatorname{div} \left( \mathbf{w} + \dot{\mathbf{u}} \right)$$

• Constitutive equations for the effective stresses (hypoplasticity with intergranular strain)







#### Weak form of the equations

$$\int_{\Omega} (\operatorname{div} \boldsymbol{\sigma} - \operatorname{grad} p + \varrho \, \mathbf{g} - \varrho \ddot{\mathbf{u}}) \cdot \delta \mathbf{u} \, d\Omega = 0 \quad \Longrightarrow$$

$$\int_{\Omega} \left[ \boldsymbol{\sigma} : \operatorname{grad} \delta \mathbf{u} - p \operatorname{div} \delta \mathbf{u} + \varrho \left( -\mathbf{g} + \ddot{\mathbf{u}} \right) \cdot \delta \mathbf{u} \right] d\Omega - \int_{\Gamma_t} \mathbf{t}^{tot} \cdot \delta \mathbf{u} \, d\Gamma = 0$$

$$\int_{\Omega} \left[ \frac{n}{K_f} \dot{p} + \operatorname{div} \left( \mathbf{w} + \dot{\mathbf{u}} \right) \right] \delta p \, d\Omega = 0 \quad \Longrightarrow$$

$$\int_{\Omega} \left[ \left( \frac{n}{K_f} \dot{p} + \operatorname{div} \dot{\mathbf{u}} \right) \delta p + \frac{k}{g} \left( \frac{1}{\varrho_f} \operatorname{grad} p - \mathbf{g} + \ddot{\mathbf{u}} \right) \cdot \operatorname{grad} \delta p \right] d\Omega 
- \int_{\Gamma_g} q \, \delta p \, d\Gamma = 0$$

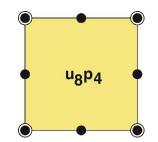




#### **Spatial discretisation**

Taylor-Hood-formulation

$$\mathbf{u} = N_I^u \mathbf{u}_I, \quad p = N_J^p p_J$$
  $\delta \mathbf{u} = N_I^u \delta \mathbf{u}_I, \quad \delta p = N_J^p \delta p_J$ 



System for the nodal variables u<sub>1</sub>, p<sub>2</sub> and their time derivatives:

$$\begin{split} &\int_{\Omega} \left( \boldsymbol{\sigma} \cdot \operatorname{grad} N_{I}^{u} - \varrho N_{I}^{u} \mathbf{g} \right) d\Omega - p_{J} \int_{\Omega} N_{J}^{p} \operatorname{grad} N_{I}^{u} \, d\Omega \\ &+ \ddot{\mathbf{u}}_{K} \int_{\Omega} \varrho N_{I}^{u} N_{K}^{u} d\Omega - \int_{\Gamma_{t}} \mathbf{t}^{tot} N_{I}^{u} d\Gamma = 0 \\ &\dot{p}_{L} \int_{\Omega} \frac{n}{K_{f}} N_{L}^{p} N_{J}^{p} d\Omega + \dot{\mathbf{u}}_{K} \cdot \int_{\Omega} N_{J}^{p} \operatorname{grad} N_{K}^{u} \, d\Omega \\ &+ p_{L} \int_{\Omega} \frac{k}{g \varrho_{f}} \operatorname{grad} N_{L}^{p} \cdot \operatorname{grad} N_{J}^{p} \, d\Omega - \int_{\Omega} \frac{k}{g} \, \mathbf{g} \cdot \operatorname{grad} N_{J}^{p} \, d\Omega \\ &+ \ddot{\mathbf{u}}_{I} \cdot \int_{\Omega} \frac{k}{g} N_{I}^{u} \operatorname{grad} N_{J}^{p} \, d\Omega - \int_{\Gamma_{q}} q N_{J}^{p} d\Gamma = 0 \end{split}$$







#### Time integration + solution process

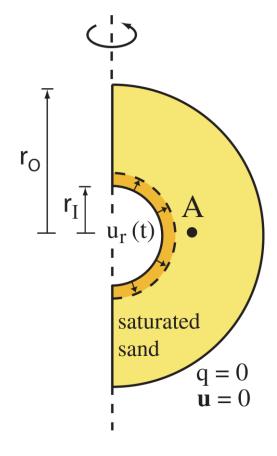
$$\begin{split} \mathbf{M}\ddot{\mathbf{d}} - \mathbf{G}\left(\mathbf{d}, \dot{\mathbf{d}}\right) &= \mathbf{0} \qquad \text{with} \quad \mathbf{d} = \left\{\mathbf{u}_{I}, p_{J}\right\} \\ \downarrow \\ \mathbf{M}\ddot{\mathbf{d}}_{(t+\Delta t)} - \mathbf{G}\left(\mathbf{d}_{(t+\Delta t)}, \dot{\mathbf{d}}_{(t+\Delta t)}\right) &= \mathbf{F}\left(\mathbf{d}_{(t+\Delta t)}, \dot{\mathbf{d}}_{(t+\Delta t)}, \ddot{\mathbf{d}}_{(t+\Delta t)}\right) = \mathbf{0} \\ \downarrow & \text{HILBER} - \text{HUGHES} - \text{TAYLOR} \\ \ddot{\mathbf{d}}_{(t+\Delta t)} &= \mathbf{f}_{1}\left(\mathbf{d}_{(t+\Delta t)}, \mathbf{d}_{(t)}, \dot{\mathbf{d}}_{(t)}, \ddot{\mathbf{d}}_{(t)}\right) \\ \downarrow & \dot{\mathbf{d}}_{(t+\Delta t)} &= \mathbf{f}_{2}\left(\mathbf{d}_{(t+\Delta t)}, \mathbf{d}_{(t)}, \dot{\mathbf{d}}_{(t)}, \ddot{\mathbf{d}}_{(t)}\right) \\ \downarrow & \dot{\mathbf{d}}_{(t+\Delta t)} &= \mathbf{F}\left(\mathbf{d}_{(t+\Delta t)}, \mathbf{d}_{(t)}, \dot{\mathbf{d}}_{(t)}, \ddot{\mathbf{d}}_{(t)}\right) \\ \downarrow & \mathbf{F}\left(\mathbf{d}_{(t+\Delta t)}, \mathbf{d}_{(t)}, \dot{\mathbf{d}}_{(t)}, \ddot{\mathbf{d}}_{(t)}\right) &= \mathbf{F}\left(\mathbf{d}_{(t+\Delta t)}\right) \overset{!}{=} \mathbf{0} \\ \downarrow & \mathbf{NEWTON} - \mathbf{RAPHSON} \\ \downarrow & \mathbf{AMATRX} & \mathbf{RHS} \\ \downarrow & \mathbf{d}^{(i+1)} &= \mathbf{d}^{(i)} + \mathbf{c}^{(i+1)} & \begin{bmatrix} \frac{\partial \mathbf{F}}{\partial \mathbf{d}} \end{bmatrix}^{(i)} \cdot \left\{\mathbf{c}\right\}^{(i+1)} &= \left\{-\mathbf{F}\right\}^{(i)} \end{split}$$

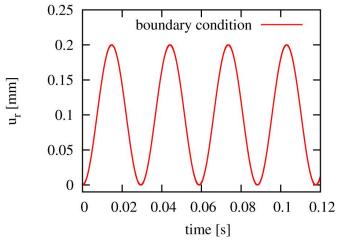




#### Verification of the user element

Spherically symmetric problem





 Verification by the comparison of the finiteelement solution (user element) with the 1D finite-difference solution

- Inner radius r<sub>i</sub>=15 cm
- Outer radius r<sub>O</sub>=5 m
- Finite element size 1.2 cm to 20 cm
- Impermeable outer boundary with zero displacement
- Vibration frequency 34 Hz

- Hypoplasticity with intergranular strain
- $K_f = 100 \text{ MPa}$
- Isotropic initial effective stress -50 kPa
- Pore pressure 50 kPa
- Void ratio 0.6
- Soil permeability *k*=10<sup>-3</sup> m/sec



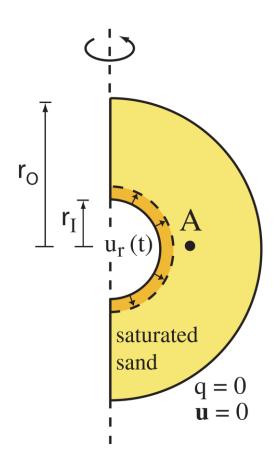


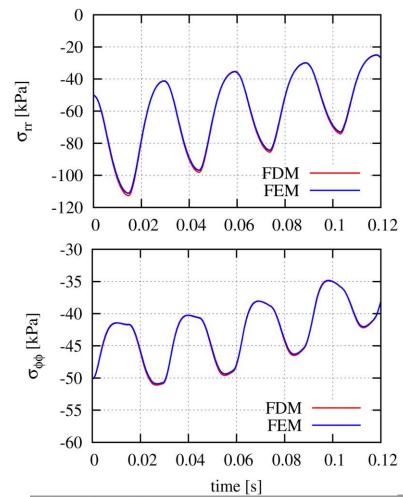


#### Verification of the user element

Comparison with 1D spherically symmetric finite-difference solution

Radial and circumferential effective stresses vs. time at Point A  $(r_A = 20 \text{ cm})$ 



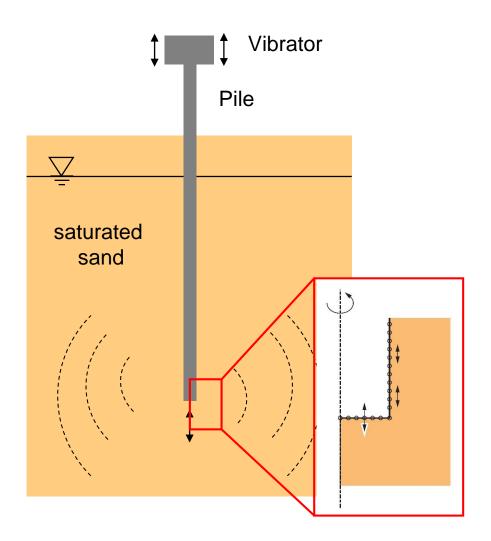








#### **Geotechnical problem**



#### Simplified Modeling technique:

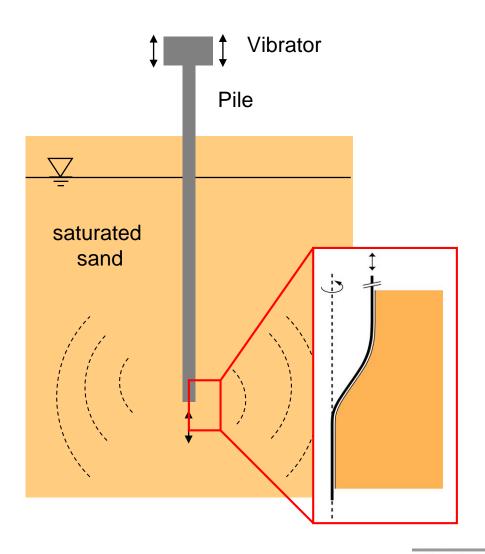
- Pile not explicitly modeled
- Pile-soil interaction considered through imposed cyclic soil displacements at the interface
- No or very low penetration







#### **Geotechnical problem**



#### Enhanced modeling technique:

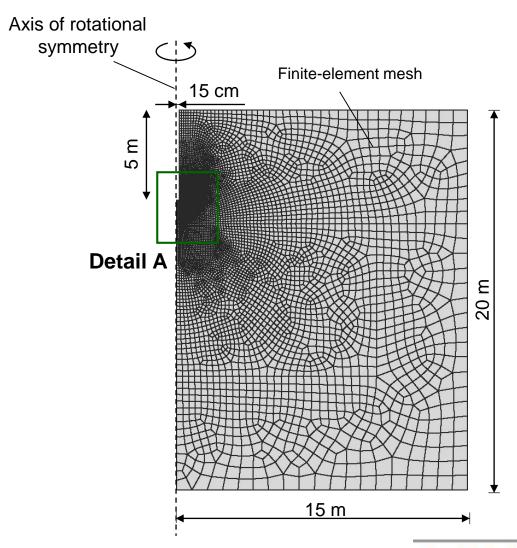
- Pile as rigid body
- Contact definition at pile-soil interface
- zipper-type technique





#### Dynamic boundary value problem (simplified model)

S. Chrisopoulos, V.A. Osinov, T. Triantafyllidis (2016): Dynamic problem for the deformation of saturated soil in the vicinity of a vibrating pile toe. Lect. Notes Appl. Comput. Mech., 80:53–67



- 2D axisymmetric formulation
- Hypoplasticity with intergranular strain (Karlsruher Sand)
- Full saturated soil (K<sub>f</sub> = 2.2 GPa)
- Geostatic initial effect. stress ( $K_0$ =1)
- Void ratio 0.6 (dense)
- Soil permeability:
  - k = 0 m/s (local undrained)
  - $k = 10^{-4}$ , 10<sup>-3</sup> m/s (local drained)
- Vibration frequency 34 Hz
- $\Delta t = 10^{-4}$  s (Period/ $\Delta t \approx 300$ )



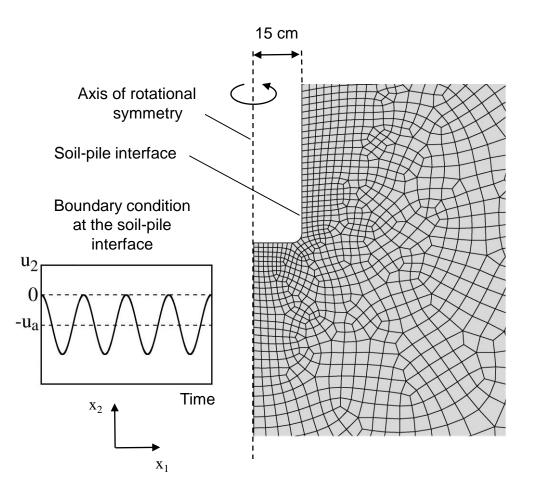




#### Dynamic boundary value problem (simplified model)

S. Chrisopoulos, V.A. Osinov, T. Triantafyllidis (2016): Dynamic problem for the deformation of saturated soil in the vicinity of a vibrating pile toe. Lect. Notes Appl. Comput. Mech., 80:53–67

#### **Detail A**



- 2D axisymmetric formulation
- Hypoplasticity with intergranular strain (Karlsruher Sand)
- Full saturated soil (K<sub>f</sub> = 2.2 GPa)
- Geostatic initial effect. stress ( $K_0$ =1)
- Void ratio 0.6 (dense)
- Soil permeability:
  - k = 0 m/s (local undrained)
  - $k = 10^{-4}, 10^{-3}$  m/s (local drained)
- Vibration frequency 34 Hz
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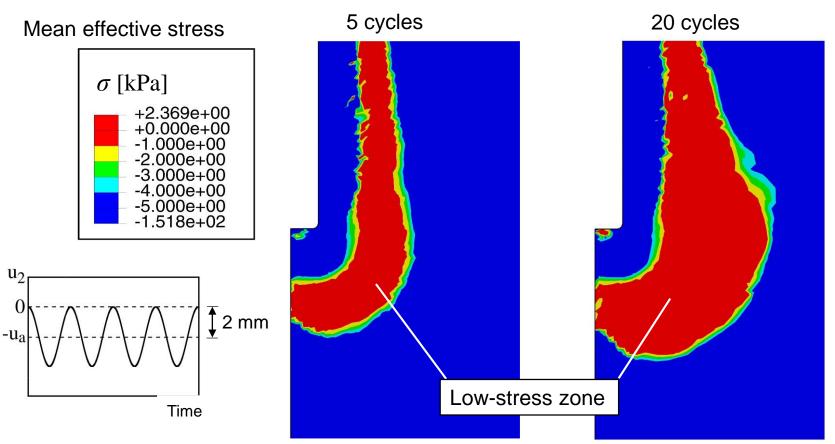






#### Numerical solution with locally undrained conditions

#### Soil permeability k = 0 m/s



#### **Result:**

 Formation of a low-stress zone (liquefaction zone) around the pile toe after several cycles of vibration

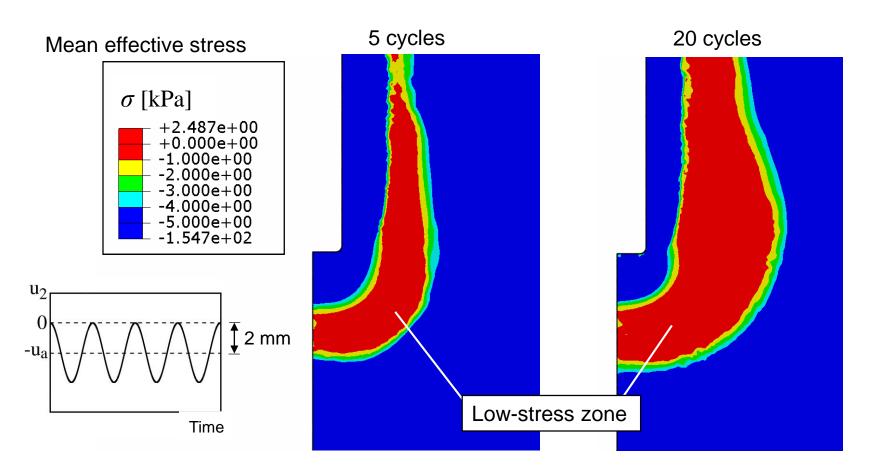






#### Numerical solution with locally drained conditions

Soil permeability  $k = 10^{-4}$  m/s

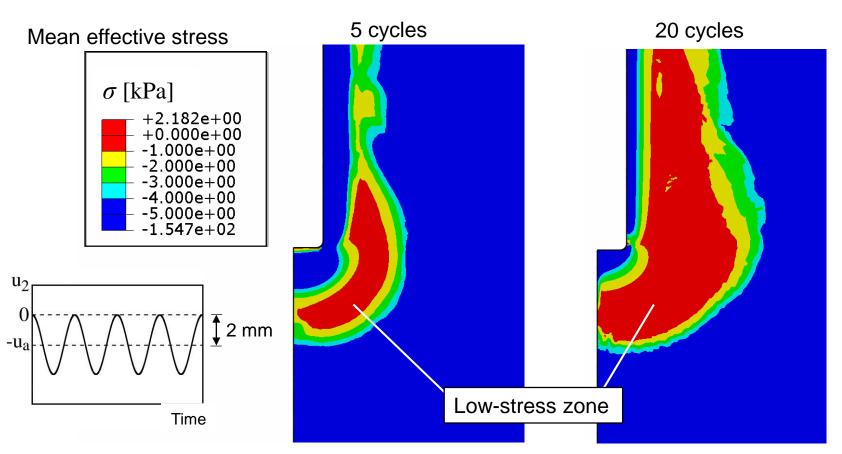






#### Numerical solution with locally drained conditions

Soil permeability  $k = 10^{-3}$  m/s



#### Result:

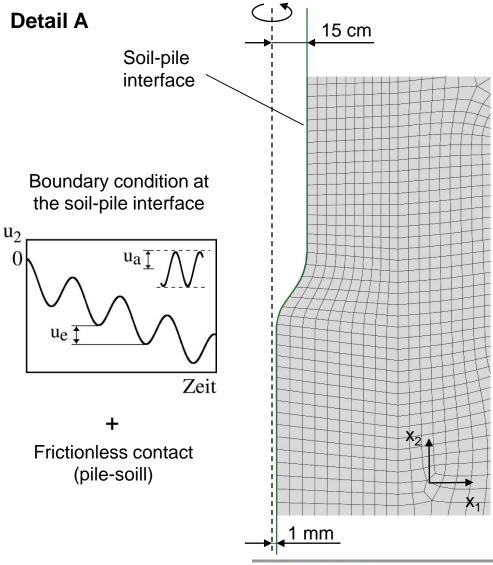
High permeability do not prevent the formation of a low-stress zone







#### Dynamic boundary value problem (enhanced model)

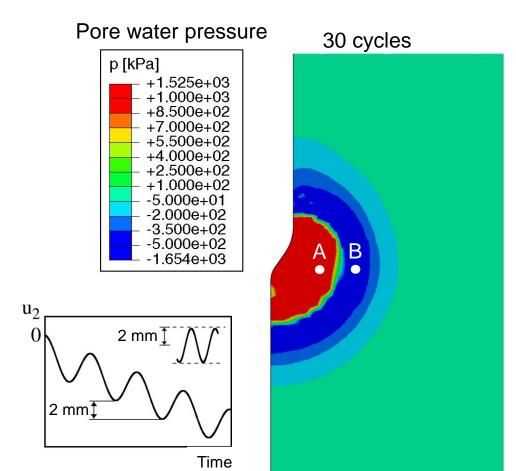


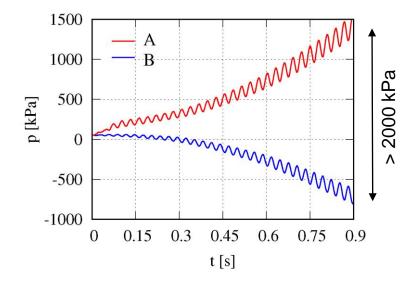




#### **Numerical solution with locally undrained conditions**

#### Soil permeability k = 0 m/s





#### Result:

 Very large, unrealistic pore water pressure gradients are created



Locally drained conditions are absolutely necessary!



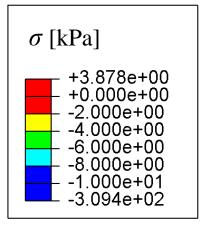


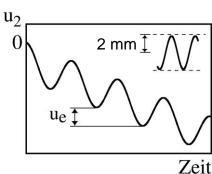


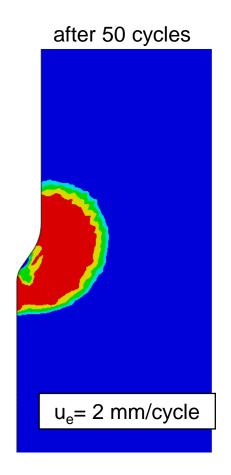
#### Numerical solution with locally drained conditions

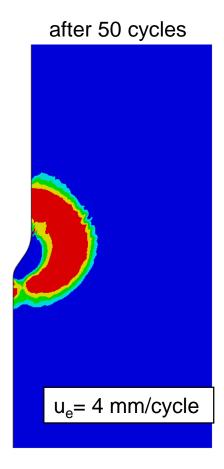
Soil permeability  $k = 10^{-3}$  m/s











#### Result:

 Despite the monotonic deformation component and the higher permeability, a lowstress zone is formed



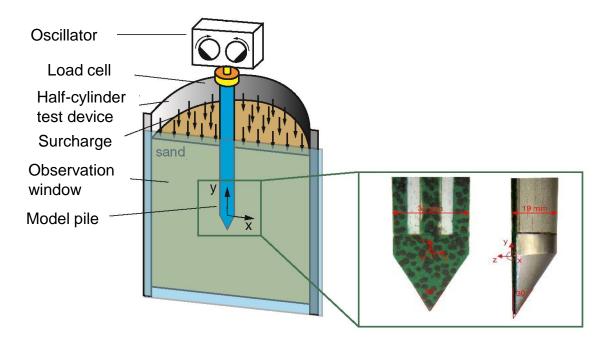




#### **Model test concept**

#### Experimental set-up (Vogelsang, 2017)

- Half-cylinder test device with observation window
- Vibration with 25 Hz for a few seconds

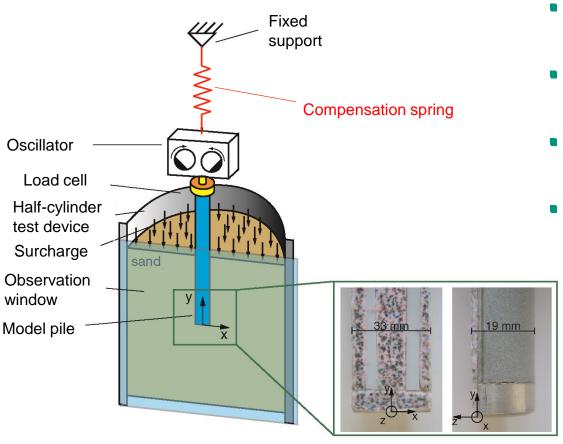






#### **Model test concept**

#### Experimental set-up (Vogelsang, 2017)



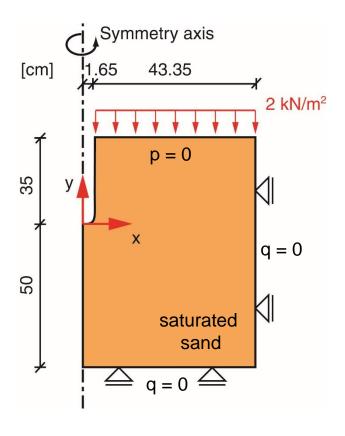
- Half-cylinder test device with observation window
- Vibration with 25 Hz for a few seconds
- Compensation spring minimizes the pile penetration
- Measurements:
  - Soil motion with DIC
  - Pile displacement and pile head force







#### Simplified FE-Model



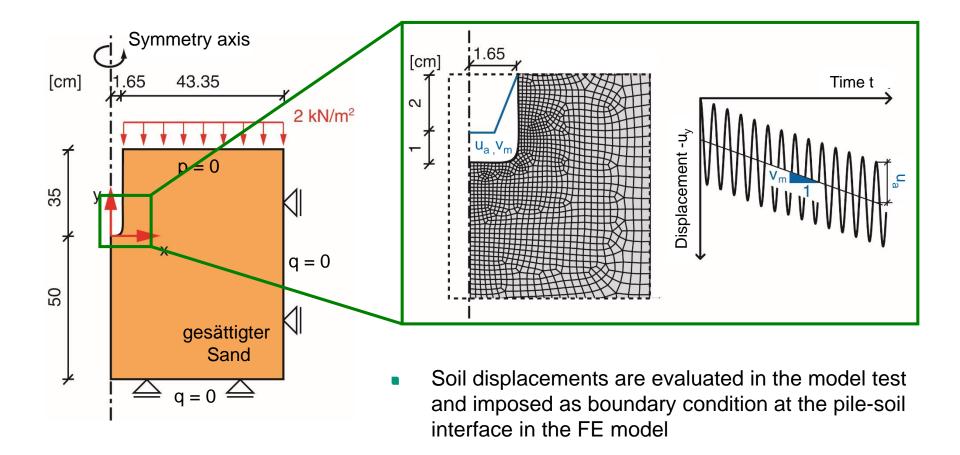
- Full saturation ( $K_f = 2.2 \text{ GPa}$ )
- Geostatic initial effective stress ( $K_0 = 0.4$ )
- Initial void ratio  $e_0 = 0.72$  (medium dense sand)
- Soil permeability:
  - $k = 1.5 \cdot 10^{-3} \text{ m/s}$  (locally drained)





#### **Simplified FE-Model**

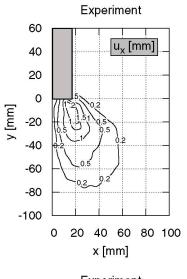
Modeling of pile-soil interaction

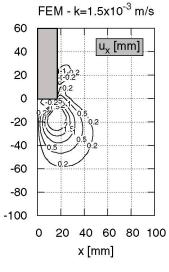


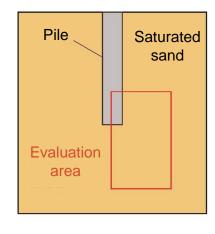


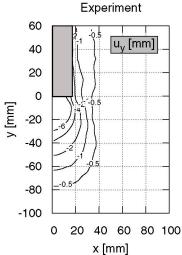


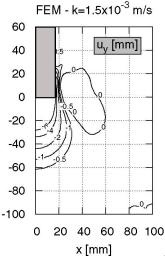
#### Displacement fields after 25 cycles











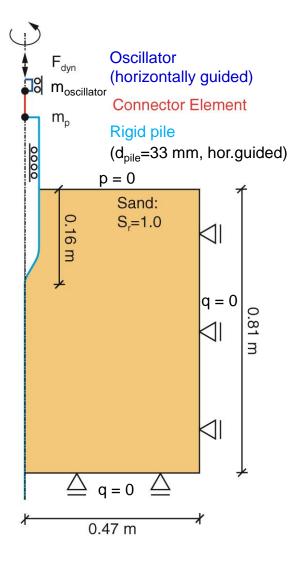
 Simplified model reproduces the soil displacements realistically







#### **Enhanced FE-Model**



- Oscillator load cell pile modeled
- Pile and oscillator rigid bodies
- Frictionless "node-to-surface" contact between pile-soil
- Initial density und soil permeability

• 
$$I_{D.0} = 0.53$$

$$k = 1.0 \cdot 10^{-3} \text{ m/s}$$

• 
$$I_{D,0} = 0.71$$

$$k = 1.2 \cdot 10^{-3} \text{ m/s}$$

$$I_{D.0} = 0.82$$

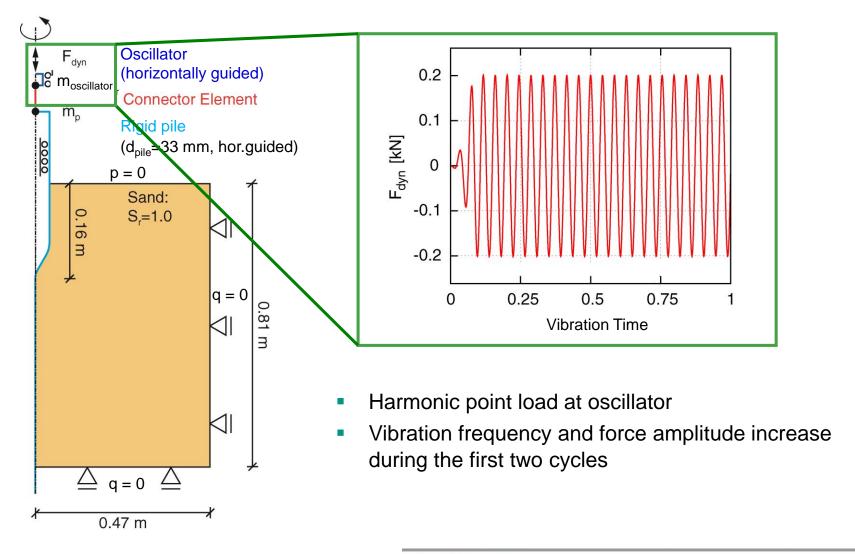
$$k = 1.4 \cdot 10^{-3} \text{ m/s}$$







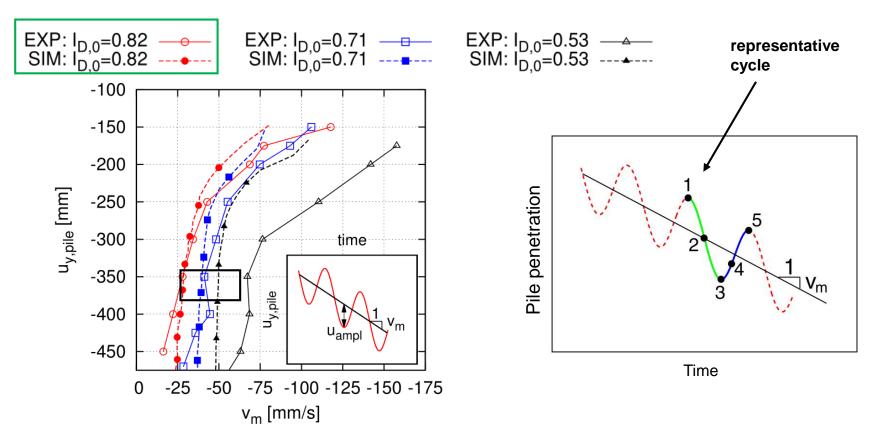
#### **Enhanced FE-Model**







Mean penetretion velocity (v<sub>m</sub>) of pile



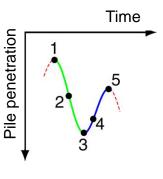
- The mean penetration velocity decreases with depth
- Comparison based on a representative cycle at a depth of about 0.4 m

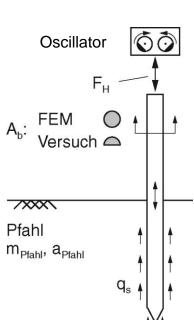




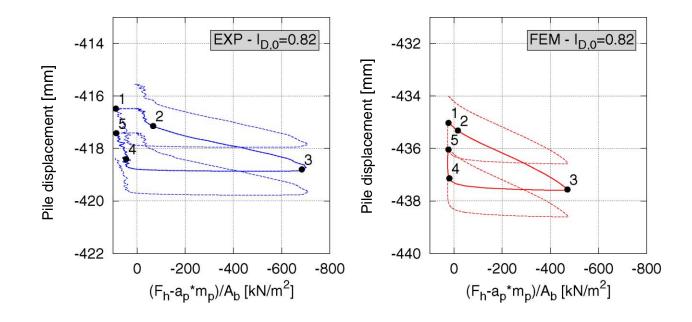


#### Related soil reaction force





 $q_b + p_w$ 



Related soil reaction force:  $\frac{F_h - m_{Pile} \cdot a_{Pile}}{A_h}$ 

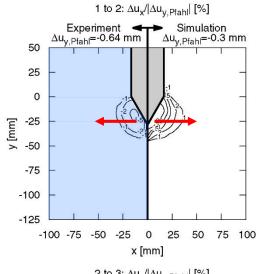
- Qualitatively similar penetration behaviour
- Lower penetration resistance in the simulation
- Shaft resistance not incorporated in the numerical model

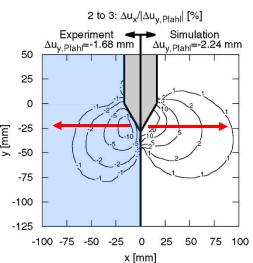


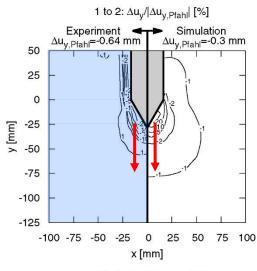


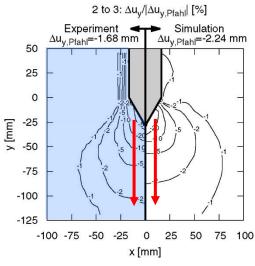


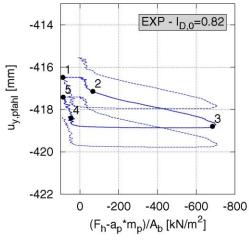
Incremental displacement fields during the penetration phase











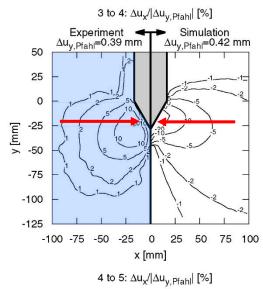
- Phase 1 ÷ 2:
- Soil displacements in a spherical area of about 1Ø under the pile tip
- The soil is directed away from the pile
- Phase 2 ÷ 3:
- Same deformation mechanism to Phase 1 ÷ 2. A deeper and larger zone of the soil is affected

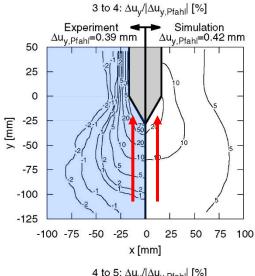


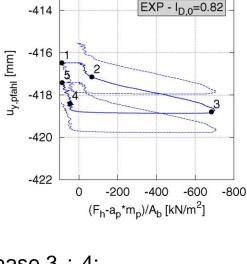


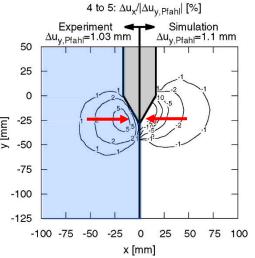


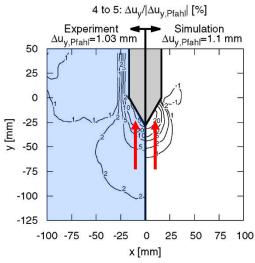
Incremental displacement fields during the upward pile motion











- Phase 3 ÷ 4:
- The soil is directed to the pile
- Phase 4 ÷ 5:
- Same deformation mechanism to Phase 3 ÷ 4
- Soil displacements in a smaller area

Chrisopoulos & Vogelsang (2019): A finite element benchmark study based on experimental modeling of vibratory pile driving in saturated sand. Soil Dyn. Earthq. Eng., 122:248-260



#### **Summary**

- A user-defined element was developed and implemented in Abaqus/Standard
- Different FE models were created for the simulation of vibratory pile driving
  - Simplified modeling approach: pile is replaced by a displacement boundary condition.
  - Enhanced modeling approach: pile as rigid body
- Vibratory pile driving in saturated soil results in a zone of nearly zero effective stresses (liquefaction zone) around the pile
- The comparison with model tests for vibratory pile driving confirms that the proposed modeling approaches are able to accurately reproduce the essential aspects of the mechanism of vibratory pile driving.



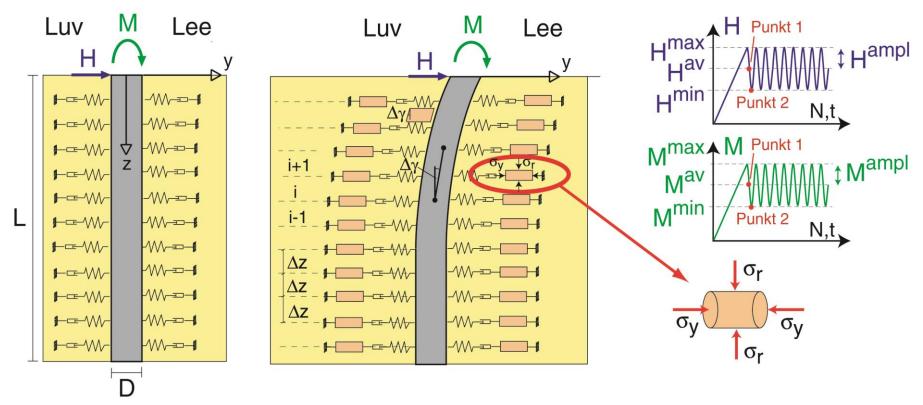


## ii. Development of an engineer-oriented model for monopile foundations based on the HCA model





Model schema with springs and dashpots arranged on both sides of the beam



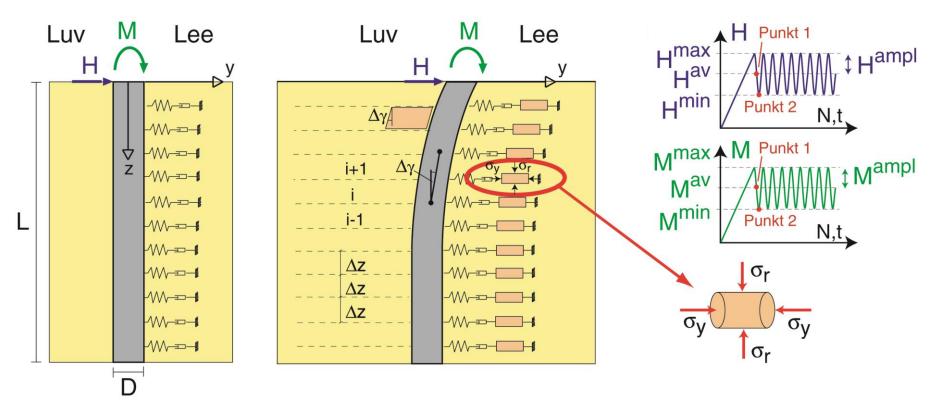
Triantafyllidis & Chrisopoulos (2016)







Model schema with springs and dashpots arranged only on one side of the beam



Triantafyllidis & Chrisopoulos (2016)





## Static loading

#### Fourth-order differential equation of the beam-on-elastic foundation problem

$$EI \cdot \frac{d^4y}{dz^4} + K_s(z) \cdot y \cdot D = 0$$

Betting modulus K<sub>s</sub> =M/D

where:  $M=((1+e_0)/C_c)\sigma_y'$  for loading

 $M = ((1 + e_0)/C_s)\sigma_y'$  for unloading

C<sub>c</sub>: compression index

C<sub>s</sub>: swelling index

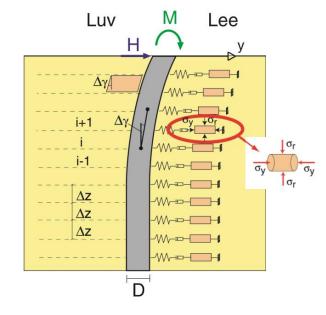
# Boundary conditions

At head: z=0

$$EI \cdot \frac{d^2y}{dz^2} = M_b$$
  $EI \cdot \frac{d^3y}{dz^3} = H$ 

At toe: z=L

$$EI \cdot \frac{d^2y}{dz^2} = 0 \qquad EI \cdot \frac{d^3y}{dz^3} = 0$$









## Cyclic loading

The deformations due to the cyclic loading are treated with equations based on the high-cycle accumulation (HCA) model. The increment of horizontal displacement  $\Delta y_i^{\rm Pile}$  of the pile is coupled with the unknown increment of stress  $\Delta \sigma_i$  in the soil caused by the cycles via:

$$[\mathbf{K}] \cdot \left[ \Delta y_i^{\text{Pile}} \right] = \left[ (-\Delta \sigma_i) \cdot \Delta z \cdot L_c \right]$$

K: stiffness matrix

 $\Delta_Z$ : distance between two springs or dashpots

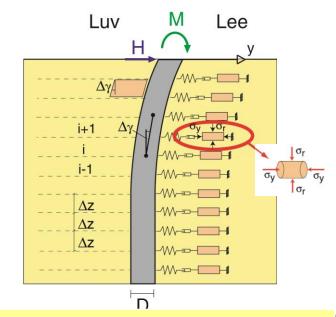
 $L_c$ =  $\pi D/4$ : characteristic length

#### **Assumption: Triaxial conditions for the soil**

From triaxial conditions:  $\Delta \sigma_i = \Delta p_i + \frac{2}{3} \Delta q_i$ 

The unknown volumetric and deviatoric stresses are approximated incrementally with the HCA model:

$$\Delta p_{i} = K(p_{i}^{0}) \left[ \Delta \varepsilon_{v_{i}} - \Delta \varepsilon_{v_{i}}^{\text{acc}} - \Delta \varepsilon_{v_{i}}^{\text{pl}} \right]$$
$$\Delta q_{i} = 3G(p_{i}^{0}) \left[ \Delta \varepsilon_{q_{i}} - \Delta \varepsilon_{q_{i}}^{\text{acc}} - \Delta \varepsilon_{q_{i}}^{\text{pl}} \right]$$



Triantafyllidis & Chrisopoulos (2016): A model for the behavior of horizontally high-cycle loaded piles. Bautechnik, 93(9): 605-627

#### Cyclic loading

$$\Delta p_i = K(p_i^0) \left[ \Delta \varepsilon_{v_i} - \Delta \varepsilon_{v_i}^{\text{acc}} - \Delta \varepsilon_{v_i}^{\text{pl}} \right] \tag{I}$$

$$\Delta q_i = 3G(p_i^0) \left[\Delta \varepsilon_{q_i} - \Delta \varepsilon_{q_i}^{acc} - \Delta \varepsilon_{q_i}^{pl}\right]$$
 (II)

Using the HCA flow rule (MCC flow rule) written in terms of Roscoe's invariants and the equation  $\Delta \varepsilon_{q_i} = \frac{2}{3} \Delta \gamma_i$  from the triaxial conditions:

$$\begin{split} \Delta \varepsilon_{v_i}^{\rm acc} &= m_v \cdot \Delta \varepsilon_i^{\rm acc}, \quad \Delta \varepsilon_{v_i}^{\rm pl} = \frac{2}{3\sqrt{3}} m_v \cdot Y \cdot \Delta \gamma_l \quad , \quad \Delta \varepsilon_{v_i} = \Delta \varepsilon_{q_i} \cdot \frac{m_v}{m_q} = \frac{2}{3} \Delta \gamma_i \quad \cdot \frac{m_v}{m_q} \\ \Delta \varepsilon_{q_i}^{\rm acc} &= m_q \cdot \Delta \varepsilon_i^{\rm acc}, \quad \Delta \varepsilon_{q_i}^{\rm pl} = \sqrt{\frac{2}{3}} \cdot Y \quad \cdot m_q \cdot \Delta \gamma_i^1, \end{split}$$

Substituting the above equations into (I) and (II) and using the indices  $\Box^0$  and  $\Box^1$  for the states at the start and the end of an increment:

$$\Delta p_i^1 = \mathcal{K}(p_i^0) \left[ \frac{2}{3} \Delta \gamma_i^1 \cdot \frac{\mathbf{m}_v^0}{\mathbf{m}_q^0} - \mathbf{m}_v^0 \cdot \Delta \varepsilon_i^{\mathrm{acc,0}} - \mathbf{m}_v^0 \cdot \frac{2}{3\sqrt{3}} \cdot \mathbf{Y}^0 \cdot \Delta \gamma_i^1 \right]$$

$$\Delta q_i^1 = 3G(p_i^0) \left[ \frac{2}{3} \Delta \gamma_i^1 - m_q^0 \cdot \Delta \varepsilon_i^{\text{acc},0} - \sqrt{\frac{2}{3}} \cdot Y^0 \cdot m_q^0 \cdot \Delta \gamma_i^1 \right]$$







#### Solution process

Using the assumption  $\Delta \gamma_i = \frac{\Delta y_{i+1}^{\text{Pile}} - \Delta y_{i-1}^{\text{Pile}}}{2\Delta z}$  the unknown volumetric and deviatoric stresses are written:

$$\Delta p_i^1 = K(p_i^0) \left[ \frac{2}{3} \cdot \frac{\Delta y_{i+1}^1 - \Delta y_{i-1}^1}{2 \Delta z} \cdot \frac{m_v^0}{m_q^0} - m_v^0 \cdot \Delta \varepsilon_i^{\text{acc,0}} - m_v^0 \cdot \frac{2}{3\sqrt{3}} \cdot Y^0 \cdot \frac{\Delta y_{i+1}^1 - \Delta y_{i-1}^1}{2 \Delta z} \right]$$

$$\Delta q_i^1 = 3G(p_i^0) \left[ \frac{2}{3} \cdot \frac{\Delta y_{i+1}^1 - \Delta y_{i-1}^1}{2 \Delta z} - m_q^0 \cdot \Delta \varepsilon_i^{\text{acc,0}} - \sqrt{\frac{2}{3}} \cdot Y^0 \cdot m_q^0 \cdot \frac{\Delta y_{i+1}^1 - \Delta y_{i-1}^1}{2 \Delta z} \right]$$

The increment of horizontal displacement  $\Delta y_i^{\rm Pile}$  of the pile is coupled with the unknown increment of stress  $\Delta \sigma_i$  in the soil via:

$$[\mathbf{K}] \cdot \left[ \Delta y_i^{\text{Pile}} \right] = \left[ (-\Delta \sigma_i) \cdot \Delta z \cdot L_c \right] \quad \text{where} \quad \Delta \sigma_i = \Delta p_i + \frac{2}{3} \Delta q_i$$





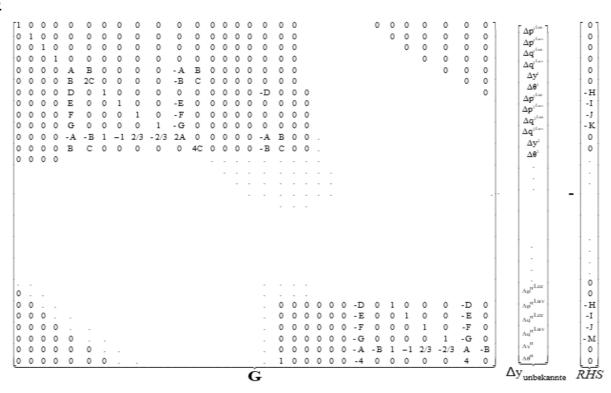


#### Solution process

#### Overall system of equations

$$[G] \cdot [\Delta y_i^{Pile}] = [RHS]$$

The system of equations is solved for each increment using the program Wolfram Mathematica



$$\begin{split} A &= \frac{12EJ}{L^3} \;,\;\; B = \frac{6EJ}{L^2} \;,\;\; C = \frac{2EJ}{L} \;,\;\; D = \frac{2}{3} \frac{m_{_{Vol}}^{\quad \text{Lee}}}{m_{_{q}}^{\quad \text{Lee}}} \frac{K(\sigma_{_{N-1}})^{\text{Lee}}}{2A_{_{Fed}}} \;,\;\; E = \frac{2}{3} \frac{m_{_{Vol}}^{\quad \text{Luv}}}{m_{_{q}}^{\quad \text{Luv}}} \frac{K(\sigma_{N-1})^{\text{Luv}}}{A_{Fed}} \;,\;\; F = \frac{2G(\sigma_{N-1})^{\text{Lee}}}{A_{Fed}} \;,\; F = \frac{2G(\sigma_{N-1})^{\text{Luv}}}{A_{Fed}} \;,\;\; F = \frac{2G(\sigma_{N-1})^{$$



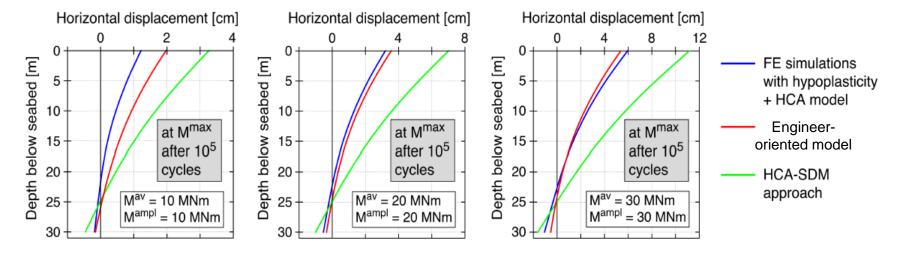




#### Results and validation

Comparison of the pile deflections of an OWPP monopile foundation after 10<sup>5</sup> cycles

<u>Parameters</u>: Pile diameter 5 m, depth of embedding 30 m, initial relative density of the soil  $I_{D0} = 0.6$  and lever arm h=M/V = 20 m



The deflection curves obtained from the engineer oriented model lie close to the results from the full FE simulations with the HCA model.

Triantafyllidis & Chrisopoulos (2016): A model for the behavior of horizontally high-cycle loaded piles. Bautechnik, 93(9): 605-627



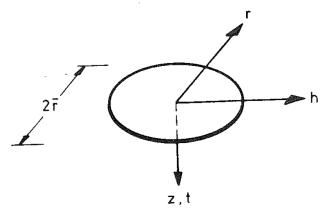
# iii. Estimation of the dynamic foundationsubsoil interaction of a wind turbine



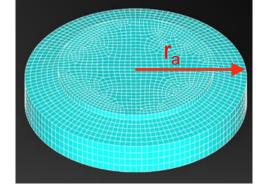


Determination of the complex dynamic stiffness of the circular foundation of the WT due to an impulse

Analytical solutions (Triantafyllidis et al. 1987)



- horizontal translation h
- vertical translation z
- rotation about r (rocking)
- rotation about z (torsion)



#### For translation:

$$Re(F_j) = u_j G r_a K_{jj}$$

$$k_{jj}$$

Im(F<sub>j</sub>) = 
$$\omega u_j r_a^2 \sqrt{G \rho} C_{jj}$$

#### For rotation:

$$Re(M_j) = \hat{u}_j G r_a^3 K_{jj}$$

Im(M<sub>j</sub>) = 
$$\omega \hat{u}_j r_a^4 \sqrt{G \rho} C_{jj}$$

K<sub>jj</sub>: Dimensionless spring parameter

C<sub>ii</sub>: Dimensionless damper parameter

r<sub>a</sub>: Radius of the circular foundation

G: Shear stiffness

 $\rho$ : Density

v<sub>s</sub>: Shear wave velocity







Calculation of the dimensionless spring K<sub>jj</sub> and damper parameter C<sub>jj</sub>

$$K_{jj} = C_0 + C_1 \cdot \alpha_0 + C_2 \cdot \alpha_0^2 + C_3 \cdot \alpha_0^3 + \dots + C_6 \cdot \alpha_0^6 = \sum_{i=0}^6 C_i \cdot \alpha_0^i$$

$$C_{ij} = C_0 + C_1 \cdot \alpha_0 + C_2 \cdot \alpha_0^2 + C_3 \cdot \alpha_0^3 + \dots + C_6 \cdot \alpha_0^6 = \sum_{i=0}^6 C_i \cdot \alpha_0^i$$

with the dimensionless frequency  $\alpha_0 = \frac{\omega r_a}{v_s}$  and the coefficients  $C_0 \dots C_6$  (from Tables)

		Kzz		Czz			
NY	0.25	0.33	0.40	0.25	0.33	0.40	
CO	-0.66796+1	-0.61402+1	-0.54920+1	-0.52469+1	-0.48135+1	-0.44293+1	
C1	-0.11174+1	+0.83234-1	-0.48845+0	-0.42338-1	-0.36151-1	-0.26038-1	
C2	+0.42391+1	+0.30600+0	+0.10339+1	-0.15837+0	-0.12768+0	-0.98526-1	
C3	-0.42496+1	+0.23483+0	-0.22400+0				
<b>C4</b>	+0.24668+1	-0.78015-1					
C5	-0.67817+0						
C6	+0.68804-1						

Table 2. Stiffness Functions Kzz (Real Part) and Czz (Imaginary Part) of Circular Foundation for Vertical Translation z







Calculation of the dimensionless spring  $K_{jj}$  and damper parameter  $C_{jj}$ 

$$K_{jj} = C_0 + C_1 \cdot \alpha_0 + C_2 \cdot \alpha_0^2 + C_3 \cdot \alpha_0^3 + \dots + C_6 \cdot \alpha_0^6 = \sum_{i=0}^6 C_i \cdot \alpha_0^i$$

$$C_{ij} = C_0 + C_1 \cdot \alpha_0 + C_2 \cdot \alpha_0^2 + C_3 \cdot \alpha_0^3 + \dots + C_6 \cdot \alpha_0^6 = \sum_{i=0}^6 C_i \cdot \alpha_0^i$$

with the dimensionless frequency  $\alpha_0 = \frac{\omega r_a}{v_s}$  and the coefficients  $C_0 \dots C_6$  (from Tables)

	Krr			Crr		
NY	0.25	0.33	0.40	0.25	0.33	0.40
CO	-0.45303+1	-0.41889+1	-0.38228+1	-0.17387-1	-0.24865-1	-0.31811-1
C1	+0.27100+0	+0.56923+0	+0.47386+0	+0.28649+0	+0.27256+0	+0.15229+0
C2	+0.79905+0	+0.17266+0	+0.17872+0	-0.20349+1	-0.18426+1	-0.12575+1
C3	-0.36608+0	-0.45346-1	-0.48259-1	+0.18237+1	+0.16423+1	+0.91185+0
C4	+0.52855-1			-0.76877+0	-0.69643+0	-0.27172+0
C5				+0.15980+0	+0.14742+0	+0.29665-1
C6				-0.13167-1	-0.12535-1	

Table 3. Stiffness Functions Krr (Real Part) and Crr (Imaginary Part) of Circular Foundation for Rotation r (Rocking)







Calculation of the complex dynamic stiffness (kij und cij)

#### **Example (vertical translation):**

#### Parameters of soil (limestone)

$$\rho$$
 = 2,6 tn/m<sup>3</sup>, G = 1,26 \* 10<sup>4</sup> MN/m<sup>2</sup>,  $v_s$  = 2200 m/s,  $r_a$  = 9,8 m,  $v$  = 0,25

for 
$$f = 5 \text{ Hz} \longrightarrow \omega = 2\pi f = 31,42 \text{ rad/s} \longrightarrow \alpha_0 = \frac{\omega r_a}{v_s} = \frac{12,57*9,9}{2200} = 0,14$$

$$K_{zz} = \sum_{i=0}^{6} C_i \cdot \alpha_0^i = -6.67 - 1.12 \cdot 0.14 + 4.24 \cdot 0.14^2 - 4.25 \cdot 0.14^3 + \cdots + 0.069 \cdot 0.14^6 = -6.76$$

$$C_{zz} = \sum_{i=0}^{6} C_i \cdot \alpha_0^i = -5.25 - 0.042 \cdot 0.14 - 0.158 \cdot 0.14^2 = -5.26$$

$$C_{ZZ} = r_a^2 \sqrt{G \rho} C_{ZZ} = -9.8^2 \cdot \sqrt{1.26 \cdot 10^4 \cdot 0.0026} \cdot 5.26 = -2.89 \cdot 10^3 \text{ MN·s/m}$$

$$Re(F_j) = u_j G r_a K_{zz} k_{zz} Im(F_j) = \omega u_j r_a^2 \sqrt{G \rho} C_{zz} C_{zz}$$







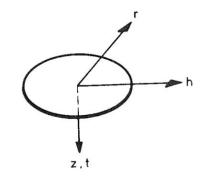
Calculation of the complex dynamic stiffness (kij und cij)

#### **Example (vertical translation):**

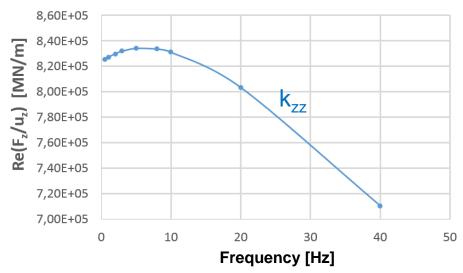
#### Parameters of soil (limestone)

$$\rho$$
 = 2,6 tn/m³, G = 1,26 \* 10<sup>4</sup> MN/m²,  $v_s$  = 2200 m/s,  $r_a$  = 9,8 m,  $v$  = 0,25

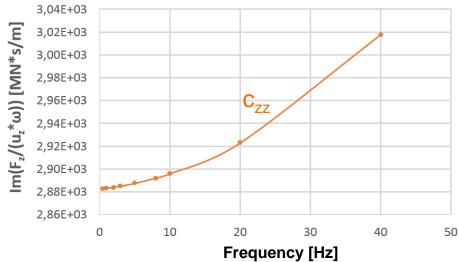
for varying frequencies:  $k_{zz}$  (f),  $c_{zz}$  (f)



#### Stiffness circular foundation (vertical translation)



#### Damping circular foundation (vertical translation)









Calculation of the complex dynamic stiffness (ki und ci)

#### **Example (rotation about r):**

#### Parameters of soil (limestone)

$$\rho$$
 = 2,6 tn/m<sup>3</sup>, G = 1,26 \* 10<sup>4</sup> MN/m<sup>2</sup>,  $v_s$  = 2200 m/s,  $r_a$  = 9,8 m,  $v$  = 0,25

for 
$$f = 5 \text{ Hz} \longrightarrow \omega = 2\pi f = 31,42 \text{ rad/s} \longrightarrow \alpha_0 = \frac{\omega r_a}{v_s} = \frac{12,57*9,9}{2200} = 0,14$$

$$K_{\hat{z}\hat{z}} = \sum_{i=0}^{6} C_i \cdot \alpha_0^i = -4.53 + 0.27 \cdot 0.14 + 0.80 \cdot 0.14^2 - 0.37 \cdot 0.14^3 + 0.052 \cdot 0.14^4 = -4.47$$

$$C_{\hat{z}\hat{z}} = \sum_{i=0}^{6} C_i \cdot \alpha_0^i = -0.017 + 0.286 \cdot 0.14 - 2.03 \cdot 0.14^2 + 1.82 \cdot 0.14^3 + \cdots -0.01 \cdot 0.14^6 = -0.012$$

$$\longrightarrow$$
  $k_{\hat{r}\hat{r}} = G r_a^3 K_{\hat{r}\hat{r}} = -1.26 \cdot 10^4 \cdot 9.8^3 \cdot 4.47 = -5.3 \cdot 10^7 MN/m$ 

$$ightharpoonup C_{rr} = r_a^4 \sqrt{G \rho} C_{rr} = -9.8^4 \cdot \sqrt{1.26 \cdot 10^4 \cdot 0.0026} \cdot 0.012 = -6.3 \cdot 10^2 \text{ MN·s/m}$$







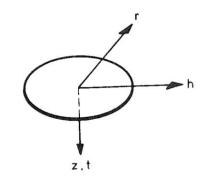
Calculation of the complex dynamic stiffness (ki und ci)

#### **Example (rotation about r):**

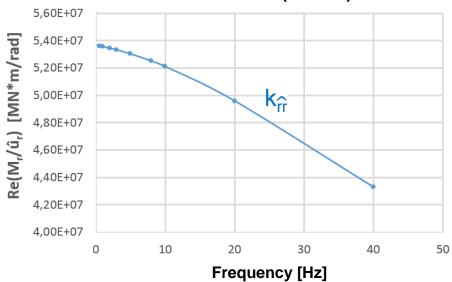
Parameters of soil (limestone)

$$\rho$$
 = 2,6 tn/m<sup>3</sup>, G = 1,26 \* 10<sup>4</sup> MN/m<sup>2</sup>,  $v_s$  = 2200 m/s,  $r_a$  = 9,8 m,  $v$  = 0,25

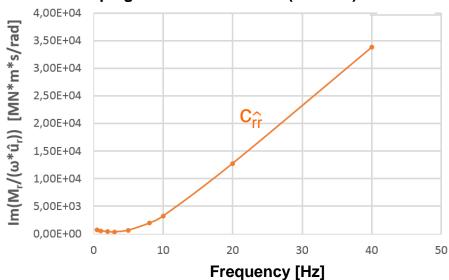
for varying frequencies:  $k_{\hat{r}\hat{r}}$  (f),  $c_{\hat{r}\hat{r}}$  (f)



#### Stiffness circular foundation (rotation)



#### Damping circular foundation (rotation)









## Thank you for your attention!





